Cloud Microphysics

Take Home Messages

Cha	apter	
3	Mol	ecular structure of water substance
Ŭ		water as the only liquid that has bigher density than its solid form
		for the vapor phase we have the ideal gas model for the solid phase the -100° C -80° C 0° C
		ideal crystal model (although coveral deviations for both exist) but for
		the liquid phase no such model eviste since it has structural properties
		the inquite prior the former two chart range order but no long range order scome to imply that H.O. melecules are
		ardered like ice in chert distance but rendemly distributed ever a long distance
		ordered like ice in short distance but randomly distributed over a long distance $mV = mPT$ or $m = -nPT$
		$pv = nRI \text{ or } p = \rho RI$
	•	the Quasi-Liquid Layer (QLL with two phases α -QLL which appears first and β -QLL always at higher
		temperatures) as the boundary between ice and other phases exists at temperatures as low as -30°C, may
		enhance chemical reactions between the ice surface and its environment, increase charge accumulation of
		graupel as well as increase ice crystal growth and multiplication
4	Bul	k thermodynamic equilibrium among water vapor, liquid water and ice
	•	if an air parcel in the atmosphere changes its equilibrium condition and the total energy is constant ($\Delta Q = 0$),
		we have adiabatic conditions ($dQ = C_v dT + p dV = 0$) and either \uparrow heat & \downarrow expansion or vice versa due to
		the 1 st law of thermodynamics when considering a closed system or an air parcel
	•	the Clausius-Clapeyron equation describes the thermodynamic equilibrium conditions of the three
		phases (assuming water vapor is an ideal gas) where the separation line between ice and liquid is slightly
		bent to the left inferring that increase in pressure during isothermal conditions leads to a phase change from
		solid to liquid form
		$\frac{dln e_{sat,w}(T)}{dln e_{sat,w}(T)} = \frac{L_e(T)}{dln e_{sat,i}(T)} = \frac{L_s(T)}{L_s(T)}$
		$dT = RT^2$ ' $dT = RT^2$
	•	the Bergeron-Findeisen process describes the transition of supercooled water into ice due to the
		environment being subsaturated with respect to supercooled water and simultaneously supersaturated with
		respect to ice and therefore ice crystals grow at the expense of supercooled liquid water
		$e_{sat,w} > e_{sat,i}$
	•	two failures and one assumption of the phase-diagram: condensation does not start as soon as
		supersaturation sets in, it is not possible to move in opposite direction of the BF-process, i.e. from ice
		crystal to supercooled water and it is assumed that materials exist in bulk quantities with no QLL
5	Sur	face thermodynamics of water substance
	•	the contact angle θ usually between 0° and 180° is a function of the liquid's properties, the substrate and the
		gas and indicates wettability and hydrophilicity of the substrate (e.g. low $\theta \rightarrow$ perfectly wettable, hydrophilic
		and many substances also hygroscopic such as sea salt)
	•	Raoult's law describes that the saturation vapor pressure over an aqueous solution droplet is decreased
		proportional to the amount of solute (in its mole fraction x _w) present due to molecular attraction by it (this is
		the solute effect) and thus a cloud droplet can survive in a subsaturated environment
		$e_{sat,w} = x_w e_{sat,w,0}$
	•	the Kelvin equation describes that the saturation vapor pressure over a cloud droplet is increased and
		is greater the smaller the radius of the curvature $(\downarrow a \triangleq \uparrow e_{sat,a})$ since molecules on a convex curved surface
		are more likely to leave it (this is the curvature effect)
		$(2v_w\sigma)$
		$e_{sat,a} = e_{sat,\infty} \exp\left(\frac{1}{RTa}\right) \rightarrow e_{sat,a} > e_{sat,\infty}$
	•	the Köhler equation here in its simplified form (with assumptions $Na^+Cl^- \rightarrow v = 2$ as well as $e_{sat a}/e_{sat co} \approx$
		1) and associated curves describe the equilibrium condition between an aqueous solution droplet and
		humid air and consists of both the Kelvin (i.e. curvature, $\pm A/a$) and Raoult (i.e. solute, $1 - B/a^3$) terms with
		the critical radius a _c and the activation supersaturation S _{act} describing the condition when the drop
		becomes activated (less supersaturation needed for bigger particles with more mass which act as better CCN)
		$S = -\frac{e_{sat,a}}{1} - 1 + \frac{A}{2} - \frac{B}{2}$ $a = -\frac{3B}{2}$ $S = -1 + -\frac{4A^3}{2}$
		$S_{v,w} = e_{sat,\infty} = 1 + a + a^{3}, \qquad u_c = \sqrt{A}, \qquad S_{act} = 1 + \sqrt{27B}$
	•	the Köhler curves can be separated into stable and unstable regimes whereas the stable regime describes
	-	the droplet growth or evaporation to a new equilibrium and the unstable part states that as soon as the
		environment is minimally supersaturated with respect to the droplet it will grow will still be in supersaturated
		conditions and thus continue its growth path to theoretical infinity
	•	S_{vw} is lower for NaCl than for (NH ₄) ₂ SO ₄ when both particles have same mass (i.e. sea salt requires less
	-	supersaturation for it to become activated) and S _{ext} is lower for binner particle masses
	1	ouporouter addition for it to boothing additionally, and water to romer for bigger particle masses

7	Nuc	cleation
	•	$\Delta \mathbf{F}_{i,}$, the free energy of formation of an i-mer , must be overcome to form a homogeneous germ and as supersaturation increases, both the germ size and the free energy required decreases $\uparrow S_{nw} \triangleq \downarrow a_{n} \triangleq \downarrow \Delta F_{n}$
	•	the population of embryos N _i follows the Boltzmann distribution with the number concentration of droplets, i.e. the population, decreasing with increasing ΔF_i , i.e. the higher the required free energy to form an i-mer, the smaller the i-mer concentration
		$\downarrow N_i \land \uparrow \land F_i$
	•	at supersaturation $S_{v,w} = 1.0$ the curve does not have a maximum which indicates that there can never be a stable size and that ΔF_i is infinite, i.e. homogeneous nucleation at RH _w = 100% is impossible
	•	homogeneous nucleation of water droplets requires a supersaturation $RH_w \sim 400\%$ and is not possible in our atmosphere and therefore heterogeneous nucleation, i.e. formation of a new phase via another particle, is mostly responsible
	•	 homogeneous ice nucleation (requires T < -38°C as well as RH_i > 140% → e.g. in high level Cirrus) starts at the subsurface when long-lasting hydrogen bonds, i.e. O····H, are formed, due to two reasons a) the freezing is accompanied by an increase in volume and the surface can accommodate that better than the aqueous bulk and b) the presence of a large electric field at the air-water interface may belo start nucleation by arranging
		molecules in preferable order
	•	tour different molecular processes can occur during heterogeneous nucleation of a liquid water drop a) desorption \rightarrow leaving embryo (also during IN _{hom.})
		b) deposition \rightarrow landing and becoming part of embryo (also during IN _{hom.})
		d) surface diffusion \rightarrow landing on substrate, diffusing randomly & some become part of embryo
	•	the size of an aerosol particle is far more important than its chemical composition to act as CCN as is
		shown by linear regression of observed vs. fixed chemical composition / fixed size distribution of CCNs criteria for effective ice nucleating particles (INP) include (see also zel)
		• water insolubility \rightarrow wrong!
		• size \rightarrow large particles (r _p > 0.1 µm) with more ice-active surface $1 \ 10^{10} \ 10^{10} \ 10^{2} \ 10^{3} \ 10^{4} \ 10^{5}$
		sites better IN CCN CRYSTALS
		• nydrophobicity \rightarrow contrasting to CCN where opposite is true • chemical bond similarity \rightarrow OH bonds in surface promote IN \rightarrow [CCN] up to 10 ² to 10 ³ higher over
		 crystallographic similarity
		• ice-active surface sites → impurities on the surface can be well as smaller droplets over land since they have to compete over LWC
		more effective than pure crystal
	•	overestimate during measurement since some crystals get scattered and break, esp. at device edges)
		i) fragmentation of ice crystals by break-up \rightarrow turbulence, collision and melting with high impact
		ii) shattering of freezing drops \rightarrow if drop freezes fast no thermal mechanical equilibrium $(p' \neq p'' \neq etc.)$
		III) Hallett-Mossop mechanism \rightarrow ejection of splinters during riming process (i.e. freezing of supercooled draplet onto crystal poods optimal temperatures (T \rightarrow 5 °C [2 °C] because if teo cold; drop freezes
		too fast without splinters & if too warm: drop spreads over surface without splinters)
7a	Kar	iji et al., 2017 – Ice formation and evolution in clouds and precipitation: measurement and modelling challenges
	٠	homogeneous freezing takes place when $T < -38^{\circ}C$ and $RH_i > 140\%$, otherwise we have heterogeneous
		nucleation alded by aerosol particles (INPs) which lower the energy barrier (ΔF_i) for germ formation and may also do so at T < -38°C with nucleation rates competitive enough when RH < 140%
	•	types of atmospheric relevant INPs include
		1.) mineral and desert dusts → most important INP – particle size more important than chemical composition – large emission rates – Microcline K-feldspar, i.e. KAISi ₃ O ₈ , most active ice component
		2.) metals and metal oxides \rightarrow efficiency varies with chemical composition – smelting, aircraft engines, etc.
		3.) organics and glassy particles \rightarrow IN ability of SOA highly varied – onset of ice nucleation is a function of temperature – pure organics promote IN via deposition mode only at T < 40°C
		4.) bioaerosols \rightarrow airborne biological aerosols such as bacteria, fungal spores, virus, etc. – as \downarrow T, more
		sites on bioparticle become active - can also be attached to other particles and thus enhance IN ability
		5.) soil dust \rightarrow excellent IN ability due to internal mixing with organic matter – 25% of global dust emissions
		o.) piomass compustion aerosols \rightarrow i.e. soot & asn with high emission rates and high atm. concentration – the IN ability appears to be linked to fuel source & mixing state – not found in Ci clouds
		7.) volcanic ash & soot particles \rightarrow act as good INPs if immersed in liquid droplets via deposition mode
		8.) crystalline salts \rightarrow NaCl recognized as good CCN but also ~25% sea salt detected in ice particles –
	_	(INH ₄) ₂ SU ₄ good candidate for heterogeneous freezing in Ci
	•	Tour modes exist for neterogeneous nucleation

		 a) deposition ice nucleation → from vapor directly to ice phase where pores help to create ice germs b) contact freezing → occurs either from 'inside-out' or from the outside when INP collides with an airwater interface
		c) immersion freezing \rightarrow INP immersed in supercooled liquid and then freezes – however the particle does not set as a CCN before publication
		d) condensation freezing \rightarrow the IN initially serves as CCN to form a drop at T < 0°C – not established yet
	•	if it is truly different from a) or c) known pre-requisites of ice-nucleating particles
		• for insoluble/solid particles (e.g. mineral dusts) \rightarrow IN natural sources anthropogenic sources
		occurs at 'ice-active surface sites ' – small particles act less likely as INP – OH-bonds at surface, lattice
		structure, size, surface defects and nature of chemical - ocean - deforestation
		• for <u>soluble particles</u> (e.g. biological ice nucleators) \rightarrow
		may separate from their carriers – water molecules \rightarrow materials in urban pollution may act to deactivate
		active sites – as temperature decreases, more sites on
		bioparticle become active – those sites bound to host particle such as mineral dust and thus provide active sites or chemical bond necessary for stabilizing IN
		 ammonium sulphate (NH₄)₂SO₄) and other solid soluble salts can form ice heterogeneously (see
9	Diff	Kanji et al., 2017 Figure 1-8 \rightarrow sea sait (NaCi) detected in ~25% of Cirrus cloud ice particle residuals usion growth and evaporation of cloud and precipitation particles
	•	in order to have diffusion growth we need to have a supersaturated environment, i.e. the water vapor
		density in the surrounding air is higher than just above the droplet surface $\rho_{v,\omega} > \rho_{v,q}$
	•	in the temperature range -40°C to 40°C Hall and Prusppbacher (1976) gave an empirical relation for the
		temperature and decreasing pressure
		$\uparrow D_v$ when $\uparrow T$ and $\downarrow p$
		jumping length close to the drop surface, we know that the mass growth rate, i.e. dm/dt , is greater higher
		up in the cloud, the larger the drop, the greater its mass growth rate and the smaller the drop, the greater its linear growth rate (at 40 µm asymptotically slow linear growth rate but it may continue under favourable
		conditions, broaden the drop size spectrum and thus help in rain formation which is otherwise not possible
		without considering the effect of collision-coalescence $dm = da = D_n^* \rho_{n,m} - \rho_{n,n}$
		$\frac{1}{dt} = 4\pi D_v^* a (\rho_{v,\infty} - \rho_{v,a}) \qquad \rightarrow \qquad \frac{1}{dt} = \frac{1}{a} \frac{1}{\rho_w} \rho_w$
	•	diffusional growth also implies a phase change from the vapor to the liquid phase and incurs the release of latent heat which warms the drop's surface and thus it is this surface temperature that determines saturation
	•	vapor density p _{v,a} at the surface ventilation due to falling motion of droplet effect enhances both diffusion growth & evaporation rates
		which can be 10x higher than for a stationary drop and is the same if p, T & N _{Re} ($N_{Re} = Ud/v$) are the same
		upstream & a reduced gradient in the wake
	•	for the diffusion growth of ice crystals , due to their many different forms, we define the capacitance C , a guantity regardless of the shape of the conductor and derive the diffusional growth of said conductor
	•	the capacitance parameter poses some problems since hollow and solid crystals (where hollow ones grow up
	•	each temperature / supersaturation ratio has a distinct ice crystal formation habit and habits may change
10	Cal	during growth with varying T & RH _i , i.e. at different conditions, growth occurs along different axes
10	•	collision growth of small drops is inefficient (the small a ₂ due to low inertia follow the streamlines of the
		bigger droplets \rightarrow small E) – becomes more efficient with turbulence (\uparrow ~25%) and when the drops are large and of approx, equal sizes (higher inertia and more deflection from the streamlines of the collector droplet)
	•	coalescence efficiency is low for large colliding pairs but is higher for large drops collecting small droplets
	•	both graupel (original crystal completely masked by rimes) and hail (graupel > 5 mm) can go through dry and wet growth regimes with the Schuman-Ludian limit defined as the boundary between the two
		regimes showing dry growth layers consisting of small polycrystalline grains and wet growth layers with large
	•	the larger the crystal, the higher the collision efficiency E _{coll} with a droplet, as long as drop is not too big
		and falls faster than crystal

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	•	riming preferentially takes place near the edge of the plate and almost completely on down-facing side
	•	the collision efficiency of ice crystals is higher with higher temperature with the maximum close to 0°C due to near-surface melting and the influence of the QLL which increases sintering , i.e. sticking together the four breakup types, with lower pressure increasing fragment numbers are as follows
		a) filament breakup \rightarrow small drop hits large one near equator – separation due inertia with filament b) sheet breakup \rightarrow small drop hits close to center – rips whole sheet out that gets unstable & desintegrates
		c) disk breakup \rightarrow small drop hits big one near center – their conclomerate forms disk which breaks up d) bag breakup \rightarrow when a disk is inflated by air – forms a bag which then desintegrates into many fragments
11	Clo	ud drop population dynamics in the warm rain process
•••		magnitude of undraft influences cloud drop size with larger drop sizes or colloidally stable cloud when:
		hydrometeors when undraft is higger
		stochastic collection process is unequal growth rates more efficient way to
		produce precipitation-sized particles than with the continuous growth model since drop size spectra broadens faster with time and with the inclusion of diffusion growth (possible after 40 µm under right conditions) even more so
		initial cloud drop size distribution i.e. how liquid water content (LWC) is
		distributed, key in determining stability and sensitivity of rain formation in both mean drop size and spread of spectrum
	•	maritime clouds with fewer CCN and therefore larger drops are more colloidally unstable than continental clouds with same LWC and more ready to rain
	•	ultragiant nuclei (r > 10 μm), which exist in both maritime and continental clouds, and mixing can accelerate raindrop formation by collision of giant dry particle and large drop
	•	rain drops get unstable in the range of 4.3 to 6.2 mm due to induced circulation & collision with smaller drops
14	Clo	ud electricity
	•	electric charges in the atmosphere originate from two sources: a) radioactive emanations from the Earth's interior and b) cosmic rays which both form small ions with high mobility (H ⁺ and e ⁻) which attach to neutral molecules or large ions which are either simple or small ions attached to aerosol particles which have a lower mobility
	•	the atmosphere between surface and ionosphere (where plasma, i.e. an electric fluid that is neutral but reacts to external electric and magnetic fields and which increases its electric conductivity with higher ion concentration, is present) acts as an electric cavity \rightarrow however, it does not insulate entirely and thus ions move from one conducting layer to the other
	•	the Earth's surface is overall negatively charged due to the global fair-weather charge and the leakage current tries to offset this distribution by bringing negative charges back up into the atmosphere but fails to neutralize fair-weather electric field due to lightnings
	•	the electric charge distribution in thunderstorms is usually one where positive charges are located at the cloud top and negative charges in lower levels with a small positive ly charged area at cloud base (this might be due to + charged graupel which when falling down melt, release tiny - charged mini-droplets and thus remaining graupel acts positive charges)
	•	lightning and thunder processes transport negative charges to the surface to replenish the fair-weather leakage via four steps and produce shock-waves at supersonic speed
		i) stepped leader \rightarrow with freshly ionized plasma at the tip $-\vec{v} = 10^5 m s^{-1}$
		 ii) streamer → going up from the ground to meet head of leader & establish continuous channel iii) return stroke → what we call lightning – while electric current and light go upwards, - charges transported to the surface – rapid heating of channel creates shockwaves and thunder iv) dart leader and 2nd return stroke → return stroke can occur multiple times – impression of blinking
	•	cloud electrification (+ charges above and - below) occurs mostly over land due to higher updraft speeds than over the ocean ($\uparrow w \triangleq \uparrow CAPE$) and the lightning frequency in a cloud is a strong function of the updraft speed , but also cloud radius and the ice volume ($f = Rw^6 \overline{V_{ice}}$) \rightarrow to produce frequent lightning, large quantities of these parameters are favourable.
	٠	charging can either be inductive (i.e. occurring due to external influence of an electric field) or non-inductive (i.e. without an external field) and includes theories like
		 convective charging → + charges carried up by updraft forming a positive charge in the interior – cloud top and edges are downdraft regions & hence - drops would be carried to lower levels
		 selective ion capture (ind.) → underside of drop is + polarized and catches dominantly - charges and acquires overall - charge – as drop is falling, + ions have to catch up at top of droplet and therefore drop remains negatively charged (quantitatively unrealistic)
		 charging due to rebound of particles during collision (ind.) → transfer of charge during rebound where + charges in lower half of large particle are neutralized – after rebound large particle carries - charge & gravitational separation would result in dipole

1	10 ¹	10 ²	10 ³	10 ⁴	10 ⁵
INP			CCN	CF	RYSTALS
\langle	due to criteria		$\hat{\Gamma}$		measuremer overestimation
• Water insolubility → w • size (> chem. compo • hydrophobicity • chemical bond simila • crystallographic simila • ice-active surface site		imilarity e sites	→ 10 ² to 10 ³ higher ove polluted areas where w also have smaller drops	due to the three r i. fragmentation <i>turbuler</i> ii. shattering of iii. Hallett-Mosso	multiplication process n of crystals by break-unce, collision, melting freezing drops: $p' \neq p'' \neq etc$. pp mechanism: $-5^{\circ}C [-8^{\circ}C, -3^{\circ}C]$
cold cloud region gravitation & sed (i.e. dry and wet	s → imentation scavenging)				

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